

IV El Modelo Gaussiano

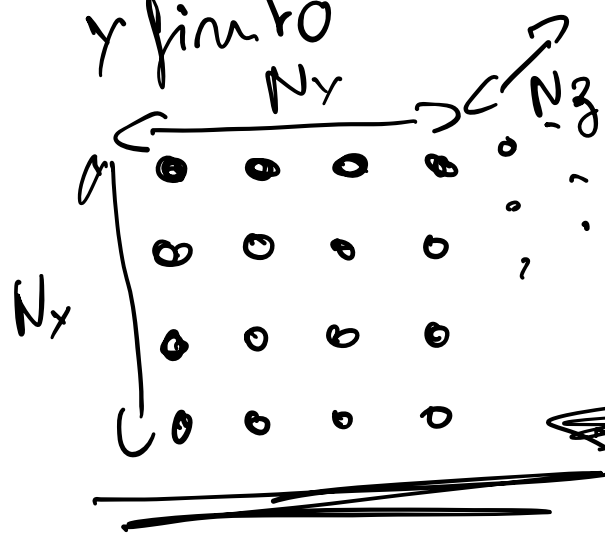
$$Z = \int \mathcal{D}\phi e^{-S(\phi)}$$

?

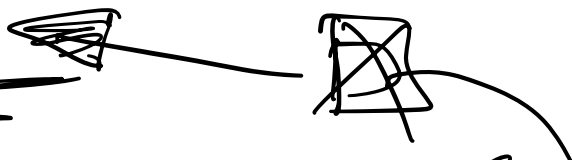
 si $S(\phi)$ es cuadrática.

Integrales múltiples en \mathbb{R}^n \rightarrow Integrales funcionales

Sistema lineal y finito

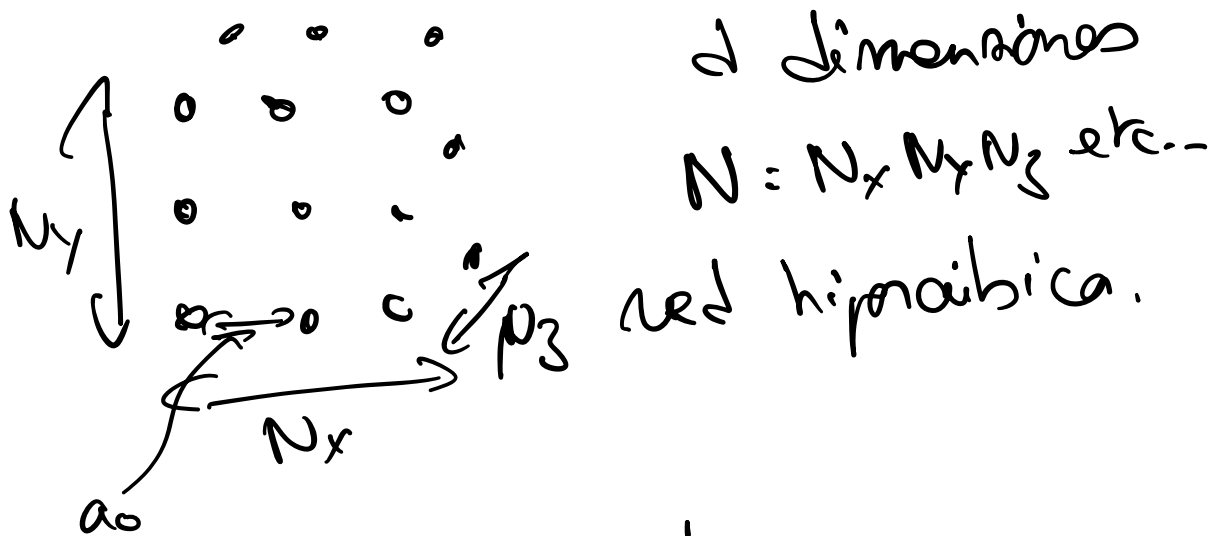


N finito
 $N = N_x \cdot N_y \cdot N_z$
 finito





Sistema discreto y finito



$$L_x = a_0 N_x$$

$$L_y = a_0 N_y$$

⋮

$$\underline{V} = L_x L_y L_z \dots$$

$$= a_0^d \underline{N}$$

$$\vec{r} = m_x a_0 \hat{e}_x + m_y a_0 \hat{e}_y + m_z a_0 \hat{e}_z + \dots$$

$m_x = 0, 1, \dots, N_x - 1$, idem for m_y etc.

Condições de borda periódicas



$$\begin{aligned} \vec{r} &\equiv \vec{r} + N_x a_0 \hat{e}_x \\ &\equiv \vec{r} + N_y a_0 \hat{e}_y \\ &\text{etc.} \end{aligned}$$

Em cada sítio, definimos

$$\phi_{\vec{r}} \in \mathbb{R} \quad \phi_{\vec{r}} = \phi_{\vec{r} + N_x a_0 \hat{e}_x}$$

Modos de Fourier:

$$e^{i\vec{k} \cdot \vec{r}} \quad \vec{r} = \frac{2\pi n_x}{L_x} \hat{e}_x + \frac{2\pi n_y}{L_y} \hat{e}_y + \dots$$

$$n_x, n_y \dots \in \mathbb{Z}$$

$$e^{i\vec{k} \cdot (\vec{r} + L_x \hat{e}_x)} = e^{i\vec{k} \cdot \vec{r}}$$

$$-\frac{N\pi}{2} < k_x \leq \frac{N\pi}{2} \quad \text{idem para } k_y, k_z.$$

~~Primer~~
Primera zona de Brillouin

⇒ ortogonalidad de los modos de Fourier:

(Todo el momento con $\forall f \neq 1$)

$$\sum_{n=0}^{N-1} \lambda^n = \frac{1 - \lambda^N}{1 - \lambda}$$

Sea muestra que:

$$\sum_{\vec{n}} \frac{e^{i(\vec{k}_1 - \vec{k}_2) \cdot \vec{n} \cdot \vec{a}}}{N} = \delta_{\vec{k}_1, \vec{k}_2}$$

$$\sum_{\vec{n} \in 2.B.} \frac{e^{i\vec{k}_1 \cdot (\vec{r}_1 - \vec{r}_2)}}{N} = \delta_{\vec{k}_1, \vec{k}_2} = \delta_{n_{x1}, n_{x2}} \delta_{n_{y1}, n_{y2}} \delta_{n_{z1}, n_{z2}}$$

Transformada de Fourier:

$$\hat{\phi}_{\vec{k}} = \int_{\mathbb{R}^3} e^{i\vec{k}\cdot\vec{r}} \phi_{\vec{r}} d\vec{r}$$

i) $\phi_{\vec{r}} = \sum_{\vec{k} \in \Gamma_B} \frac{e^{-i\vec{k}\cdot\vec{r}}}{N} \hat{\phi}_{\vec{k}}$

ii) $\phi_{\vec{r}} \in \mathbb{R} \Rightarrow \hat{\phi}_{\vec{k}}^* = \hat{\phi}_{-\vec{k}}$

iii) $\int_{\mathbb{R}^3} \phi_{\vec{r}}^2 d\vec{r} = \sum_{\vec{k}} \frac{1}{N} |\hat{\phi}_{\vec{k}}|^2$

iv) $\int_{\mathbb{R}^3} \phi_{\vec{r}} \phi_{\vec{r}+\vec{a}_n} d\vec{r}$
 \downarrow
 $\int_{\mathbb{R}^3} \phi_{\vec{r}} \phi_{\vec{r}+\vec{a}_n} d\vec{r} = \frac{1}{2} \sum_{\vec{k}} [\hat{\phi}_{\vec{k}-\vec{a}_n} \hat{\phi}_{\vec{k}} + \hat{\phi}_{\vec{k}} \hat{\phi}_{\vec{k}+\vec{a}_n}]$

$$= \frac{1}{N} \sum_{\vec{k}} |\hat{\phi}_{\vec{k}}|^2 \cos(kx a)$$

imaginemos que:

$$Z = \int \prod_{\vec{n}} d\phi_{\vec{n}} e^{-S}$$

dando S es una forma cuadrática
 \rightarrow de los $\phi_{\vec{n}}$ $\leftarrow N \times N \text{ ER}$

$$S = \sum_{\vec{n}, \vec{n}'} \phi_{\vec{n}} A_{\vec{n}, \vec{n}'} \phi_{\vec{n}'} \leftarrow N$$

\uparrow inv. por traslación

$$\text{Def } A_{\vec{n}} = A_{-\vec{n}}$$

$$A_{\vec{n}} = \sum_{\vec{k}} \frac{1}{N} e^{i\vec{k} \cdot \vec{n}} A_{\vec{k}}$$

$$S = \frac{1}{2} (S + S) = \frac{1}{2N} \sum_{\vec{k}} |\hat{\phi}_{\vec{k}}|^2 [\hat{A}_{\vec{k}} + \hat{A}_{\vec{k}}^*]$$

$$Z = \int \prod_{\vec{k}} d\hat{\phi}_{\vec{k}} e^{-S}$$

↪ $\hat{\phi}_{\vec{k}}$

obs $\hat{\phi}_{\vec{k}} + \hat{\phi}_{-\vec{k}}$ no are indep. (2.s.)

$$\prod_{\vec{k}} d\hat{\phi}_{\vec{k}} \rightarrow \prod_{\vec{k}} d\hat{\phi}_{\vec{k}}^R d\hat{\phi}_{\vec{k}}^I = \prod_{\vec{k}} d\hat{\phi}_{\vec{k}}$$

N variables $\hat{\phi}_{\vec{k}}^R = \text{Re}(\hat{\phi}_{\vec{k}})$

el Jacobiano vale 1.

$$\rightarrow Z = \int \prod_{\vec{k}} d\hat{\phi}_{\vec{k}} e^{-\frac{1}{2N} \sum_{\vec{k}} |\hat{\phi}_{\vec{k}}|^2 (\hat{A}_{\vec{k}} + \hat{A}_{\vec{k}}^*)}$$

↪ $d\hat{\phi}_{\vec{k}}^I d\hat{\phi}_{\vec{k}}^R$ $\hat{\phi}_{\vec{k}}^R + \hat{\phi}_{-\vec{k}}^I$

→ N integrales desacopladas

del estilo $\int_{-\infty}^{\infty} dx e^{-\alpha x^2} = \sqrt{\frac{\pi}{\alpha}}$

$$Z = \frac{1}{h} \frac{\pi N}{2 \operatorname{Re}(\sqrt{A} \epsilon)}$$

Def: función de correlación:

$$\langle \phi_{\vec{r}_1} \phi_{\vec{r}_2} \rangle$$

$$= \frac{1}{Z} \int \frac{1}{h} \frac{\phi_{\vec{r}_1} \phi_{\vec{r}_2}}{h} \phi_{\vec{r}_1} \phi_{\vec{r}_2} e^{-S}$$

$$= \frac{1}{Z} \frac{1}{h^2} \sum_{\vec{r}_1, \vec{r}_2} e^{i\vec{h}_1 \cdot \vec{r}_1 + i\vec{h}_2 \cdot \vec{r}_2}$$

$$x \int \frac{1}{\alpha} d\vec{\phi}_r \underbrace{\vec{\phi}_r}_{\vec{\phi}_r = \begin{pmatrix} \phi_{r1} + i\phi_{r2} \\ \phi_{r1} - i\phi_{r2} \end{pmatrix}} e^{-\alpha}$$

again : $\int dx dy e^{-\alpha(x^2+y^2)} = \frac{1}{\alpha}$

$\int dx dy x y e^{-\alpha(x^2+y^2)} = 0$

$\int dx dy (x^2) e^{-\alpha(x^2+y^2)} \neq 0$

$\int dx dy e^{-\alpha(x^2+y^2)} \times \int dx dy (x^2+y^2) e^{-\alpha(x^2+y^2)} = \frac{1}{\alpha}$

Se muestra en qro, el unico caso que da $\neq 0$ es si $\vec{k}_2 = -\vec{k}_1$!

$$\mathcal{D} \langle \phi_{\vec{k}_1} \phi_{\vec{k}_2} \rangle = \frac{1}{N^2} \sum_{\vec{k}} e^{i\vec{k} \cdot (\vec{r}_1 - \vec{r}_2)} \langle \phi_{\vec{k}} \phi_{-\vec{k}} \rangle$$

$$\times \langle \phi_{\vec{k}} \phi_{-\vec{k}} \rangle = \frac{N}{2Ne(A\vec{k})}$$

$$\mathcal{D} \langle \phi_{\vec{k}_1} \phi_{\vec{k}_2} \rangle = \frac{1}{N} \sum_{\vec{k}} \frac{1}{2Ne(A\vec{k})} e^{i\vec{k} \cdot (\vec{r}_1 - \vec{r}_2)}$$

Si A acopla los ϕ 's en el mismo sitio + primeros vecinos

$$S = \sum_{\vec{r}} \left[\alpha \phi_{\vec{r}}^2 - \beta \left(\phi_{\vec{r}} \phi_{\vec{r} + a\hat{x}} + \phi_{\vec{r}} \phi_{\vec{r} + a\hat{y}} \right) \right]$$

$$S = \frac{1}{2} \sum_{\vec{k}} |\vec{\Phi}_{\vec{k}}|^2 [\alpha - \beta(\cos(k_x a_0) + \cos(k_y a_0) + \dots)]$$

$$= \frac{1}{2} \sum_{\vec{k}} \dots$$

$$Z = \prod_{\vec{k}} \frac{\pi N}{2 [\alpha - \beta (\cos(k_x a_0) + \dots)]}$$

$$\langle \phi_{\vec{n}_1} \phi_{\vec{n}_2} \rangle = \frac{1}{Z} \sum_{\vec{k}} \frac{e^{i\vec{k} \cdot (\vec{n}_1 - \vec{n}_2)}}{2 [\alpha - \beta (\cos(k_x a_0) + \dots)]}$$

2) Límite del continuo

$$a_0 \rightarrow 0, \quad L_x \text{ finitos, } N_x = \frac{L_x}{a_0} \rightarrow \infty$$

$$L_y \quad \quad \quad N_y \rightarrow \infty$$

\vec{k} siguiendo a

$$\vec{k} = \frac{2\pi n_x}{L_x} \hat{e}_x + \frac{2\pi n_y}{L_y} \hat{e}_y + \dots$$

Para $-\infty < n_x < \infty$ etc...
 $-\infty < n_y < \infty$

$$N = \frac{V}{a_0^d} \rightarrow \infty, \quad V = L_x L_y L_z \text{ finito.}$$

diccionario para pasar del espacio discreto al continuo

$$\delta_{\vec{n}_1, \vec{n}_2} \rightarrow a_0^d \delta(\vec{r}_1 - \vec{r}_2)$$

$$\sum_{\vec{k}} \rightarrow \frac{1}{a_0^d} \int d\vec{k}$$

por ejemplo

$$\sum_{\vec{k}} \frac{e^{i\vec{k} \cdot (\vec{r}_1 - \vec{r}_2)}}{N} = \delta_{\vec{r}_1, \vec{r}_2}$$

$$\rightarrow \frac{1}{V} \sum_{\vec{k}} e^{i\vec{k} \cdot (\vec{r}_1 - \vec{r}_2)} = \delta(\vec{r}_1 - \vec{r}_2)$$

γ verändert:

$$\sum_{\vec{k}} e^{i(\vec{k}_1 - \vec{k}_2) \cdot \vec{r}} = \delta_{\vec{k}_1, \vec{k}_2}$$

$$\int d\vec{k} e^{i(\vec{k}_1 - \vec{k}_2) \cdot \vec{r}} = V \delta_{\vec{k}_1, \vec{k}_2}$$

$$\phi_{\vec{k}_1} \rightarrow \boxed{\phi(\vec{r}) = \frac{1}{\Omega} \phi_{\vec{k}_1}}$$

T.f.

$$\phi_{\vec{k}} = \sum_{\vec{r}} e^{i\vec{k} \cdot \vec{r}} \phi_{\vec{r}} \rightarrow \int d\vec{r} e^{i\vec{k} \cdot \vec{r}} \phi(\vec{r})$$

$$\phi(\vec{r}) = \frac{1}{V} \sum_{\vec{k}} e^{i\vec{k} \cdot \vec{r}} \phi_{\vec{k}}$$

ans:

$$\begin{aligned} \text{i) } \sum_{\vec{k}} \phi_{\vec{k}}^2 &\rightarrow a_0^d \int d\vec{r} \phi^2(\vec{r}) \\ &= \frac{a_0^d}{V} \sum_{\vec{k}} |\phi_{\vec{k}}|^2 \end{aligned}$$

$$\text{ii) } \sum_{\vec{k}} \phi_{\vec{k}} \phi_{\vec{k} \pm a_0 \hat{e}_x} \rightarrow \phi(\vec{r} \pm a_0 \hat{e}_x)$$

$$\rightarrow a_0^d \int d\vec{r} \phi(\vec{r}) \left(\phi(\vec{r}) \left(\frac{1}{2} a_0 \partial_x \phi(\vec{r}) + \frac{a_0^2}{2} \partial_x^2 \phi(\vec{r}) + \dots \right) \right)$$

$$\text{b) } \sum_{\vec{k}} \phi_{\vec{k}} \phi_{\vec{k} + a_0 \hat{e}_x} + \phi_{\vec{k}} \phi_{\vec{k} - a_0 \hat{e}_x}$$

$$\rightarrow a_0^d \int d\vec{r} \left[\phi^2(\vec{r}) + a_0^2 \partial_x^2 \phi^2(\vec{r}) + \dots \right]$$

$$= \frac{a_0^d}{V} \sum_{\vec{k}} |\phi_{\vec{k}}|^2 \left[1 - a_0^2 k_x^2 + \dots \right]$$

$$S = \sum_{\vec{n}} \left[\alpha \phi_{\vec{n}}^2 - \beta \left(\phi_{\vec{n}} \phi_{\vec{n} + a_0 \vec{e}_1} + \dots \right) \right]$$

$$\rightarrow \int d^d \vec{r} \left[\frac{a_2}{2} \phi^2(\vec{r}) + \frac{c}{2} (\nabla \phi)^2 + \dots \right]$$

with $a_2 = 2a_0^d (\alpha - \beta)$, $c = \beta a_0^{d+2}$

$$S = \frac{1}{V} \sum_{\vec{n}} |\hat{\phi}_{\vec{n}}|^2 \left(\frac{a_2}{2} + \frac{c \vec{k}^2}{2} \right)$$

obs: $a_2, c \geq 0$

$$\langle \phi(\vec{r}_1) \phi(\vec{r}_2) \rangle = \frac{1}{a_0^{2d}} \langle \phi_{\vec{n}_1} \phi_{\vec{n}_2} \rangle$$

$$\rightarrow \frac{1}{a_0^{2d}} \left(\frac{V}{a_0^d} \right) \sum_{\vec{n}} \frac{e^{i \vec{k} \cdot (\vec{n}_1 - \vec{n}_2)}}{2(\alpha - \beta c a_0^d (k_x a_0)^2 + \dots)}$$

para $\cos(k_x a_0) = 1 - \frac{1}{2} k_x^2 a_0^2 + \dots$

$$\rightarrow \langle \phi(\vec{r}_1) \phi(\vec{r}_2) \rangle = \frac{1}{V}$$

$$\times \sum_{\vec{k}} \frac{e^{i\vec{k} \cdot (\vec{r}_1 - \vec{r}_2)}}{a_2 + c\vec{k}^2}$$

3) limite de $V \Rightarrow \infty$

$$L_x, L_y, L_z \rightarrow \infty$$

$\Rightarrow \vec{k}$ variable continue.

$$\Delta k_x \Delta k_y \Delta k_z = \frac{2\pi}{L_x} \frac{2\pi}{L_y} \frac{2\pi}{L_z}$$

$$= \frac{(2\pi)^3}{V}$$

$$\sum_{\vec{k}} \rightarrow \int \frac{d^3\vec{k}}{\Delta k_x \Delta k_y \Delta k_z} = V \int \frac{d^3\vec{k}}{(2\pi)^3}$$

$$\delta_{\vec{k}_1, \vec{k}_2} \rightarrow \frac{(2\pi)^d}{V} \delta(\vec{k}_1 - \vec{k}_2)$$

pa ejemplo

$$\sum_{\vec{k}} \frac{e^{i\vec{k} \cdot (\vec{r}_1 - \vec{r}_2)}}{V} = \delta(\vec{r}_1 - \vec{r}_2)$$

$$\hookrightarrow \int \frac{d^d k}{(2\pi)^d} e^{i\vec{k} \cdot (\vec{r}_1 - \vec{r}_2)} = \delta(\vec{r}_1 - \vec{r}_2)$$

$$\int d^d \vec{k} e^{i(\vec{k}_1 - \vec{k}_2) \cdot \vec{r}} = V \delta_{\vec{k}_1, \vec{k}_2}$$

$$\hookrightarrow \int d^d \vec{k} e^{i(\vec{r}_1 - \vec{r}_2) \cdot \vec{k}} = (2\pi)^d \delta(\vec{r}_1 - \vec{r}_2)$$

$$\hat{\phi}(\vec{k}) \equiv \int d^d \vec{r} e^{i\vec{k} \cdot \vec{r}} \phi(\vec{r})$$

$$\phi(\vec{r}) = \frac{1}{(2\pi)^d} \int d^d \vec{k} e^{-i\vec{k} \cdot \vec{r}} \hat{\phi}(\vec{k})$$

$$S = \frac{1}{V} \sum_{\vec{k}} |\hat{f}_{\vec{k}}|^2 \left(\frac{a_2}{2} + \frac{c}{2} \vec{k}^2 \right)$$

$$\rightarrow \int \frac{d^3 \vec{k}}{(2\pi)^3} |\hat{\phi}(\vec{k})|^2 \left(\frac{a_2}{2} + \frac{c}{2} \vec{k}^2 \right)$$

$$\star \langle \phi(\vec{r}_1) \phi(\vec{r}_2) \rangle = \int \frac{d^3 \vec{k}}{(2\pi)^3} \frac{e^{i\vec{k} \cdot (\vec{r}_1 - \vec{r}_2)}}{a_2 + c \vec{k}^2}$$

$$\langle \phi(\vec{r}_1) \phi(\vec{r}_2) \rangle = G(\vec{r}_1 - \vec{r}_2)$$

da da

$$G(\vec{r}) = \int \frac{d^3 \vec{k}}{(2\pi)^3} \frac{e^{i\vec{k} \cdot \vec{r}}}{\vec{k}^2 + m^2}$$

$$\vec{c} = \frac{1}{c (2\pi)^3}, \quad m^2 = \frac{a_2}{c} \geq 0$$

$$G(\tau) \xrightarrow{|\tau| \rightarrow \infty} \frac{e^{-\frac{|\tau|}{\gamma}}}{\tau \dots}$$

daíde $\gamma = \frac{1}{m}$: la longitud de correlación.