

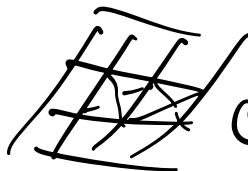
## VII El grupo de renormalización

1) En los puntos críticos  
los sistemas son invariantes  
de escala, y ademas hay

Universalidad:

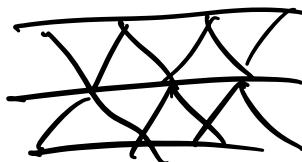
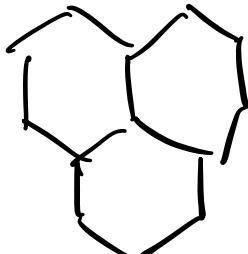
Ejemplo:

Modelo de Ising ferro



cuadrada a triangular

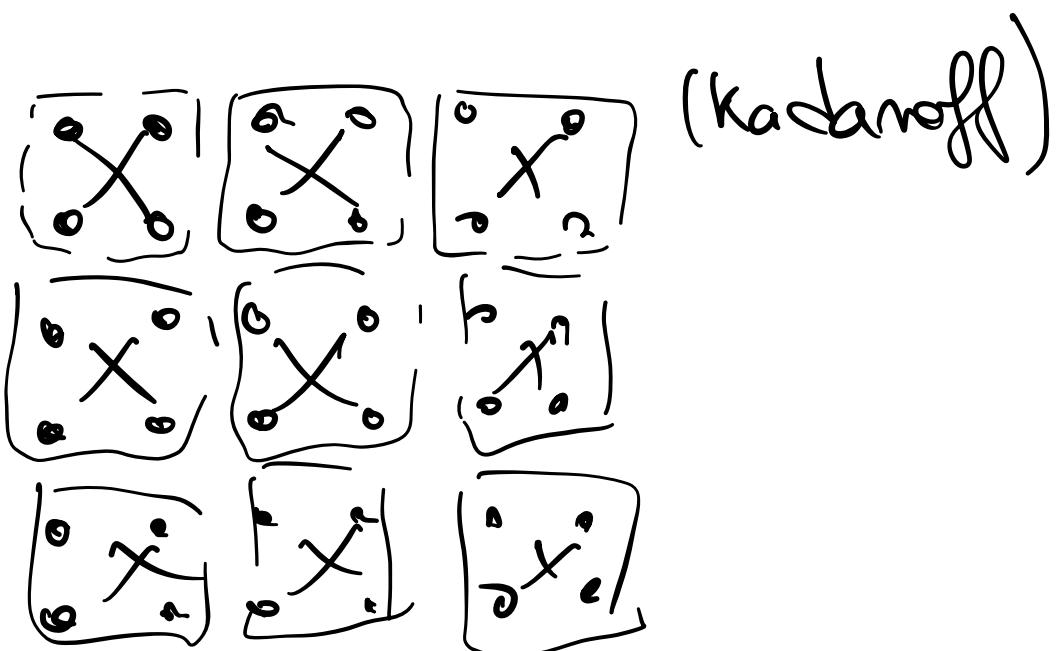
a:



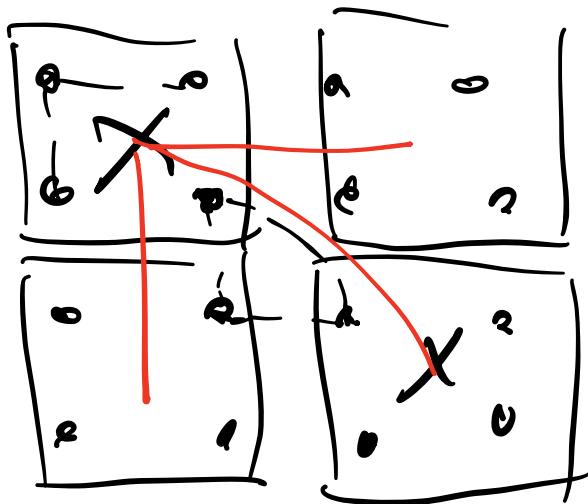
Tener  $T_c$  que sea diferente, para la clase de universalidad (los exp. artificiales) sea las mismas

Idea: "Zoom out" del modelo

2) R.C. en el espacio real



$$H(\tau, \mathfrak{F}) \longrightarrow \tilde{H}(\tilde{\tau}, \tilde{\mathfrak{F}})$$



3) El caso de  $D=1$  ( $h=0$ )

$$Z = \sum_{\{\tau_i\}} e^{k \sum_i \sigma_i \tau_{i+1}}$$

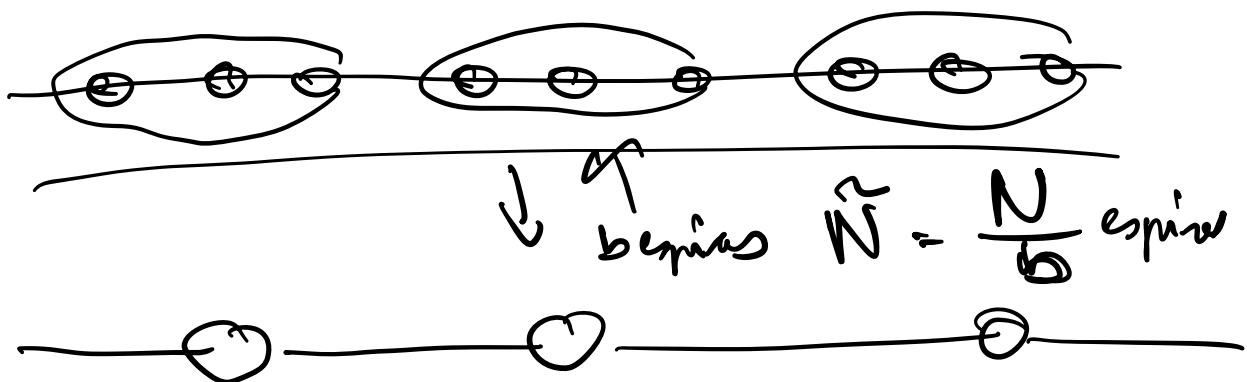
$$K = \frac{S}{k_B T}$$

$$Z = T_2 \int T^N$$

$$T = \begin{bmatrix} e^{ik} & c^{-k} \\ c^{-k} & e^{-ik} \end{bmatrix}$$

en la base  $\frac{1}{\sqrt{2}}(|\uparrow\rangle + |\downarrow\rangle)$  y  $\frac{1}{\sqrt{2}}(|\uparrow\rangle - |\downarrow\rangle)$

$$T = \begin{pmatrix} 2ch k & 0 \\ 0 & 2sh k \end{pmatrix} \xleftarrow{N \text{ copias}}$$



$$Z = T_2 \int \tilde{T}^{\tilde{N}}$$

$$\tilde{N} = \frac{N}{b}$$

$$T^N = (T^b)^{\frac{N}{b}}$$

$$\Rightarrow \hat{T} = \frac{T^b}{T}$$

$$\underline{\underline{T^b = \hat{T} = \frac{c\kappa}{\epsilon}} \left( \begin{array}{cc} 2\text{ch}\hat{x} & 0 \\ 0 & 2\text{sh}\hat{x} \end{array} \right) \rightarrow}$$

$$\text{th}(\hat{x}) = (\text{th } x)^b$$

$$\Rightarrow \hat{x} = \text{th}^{-1}[(\text{th } x)^b]$$

$$x = \frac{J}{k_B T} \Rightarrow \hat{x} = \frac{J}{k_B T}$$

$$0' = \frac{J}{k_B T}$$

$$\boxed{\hat{x} = \text{th}^{-1}[(\text{th } x)^b]}$$

extra reanaria time 2 pentos

- if  $x = 0$  ( $T = \infty$ )
- $\text{th}(x) = 1$  ( $T = 0$ )



$$T=0$$



$$T=\infty$$



$$\Delta m = \langle \sigma \rangle = 0$$

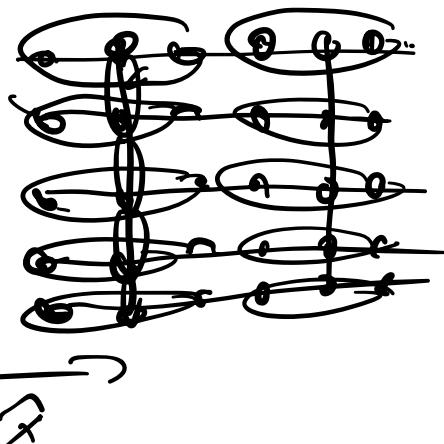
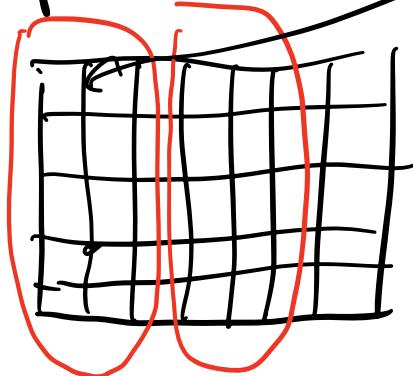


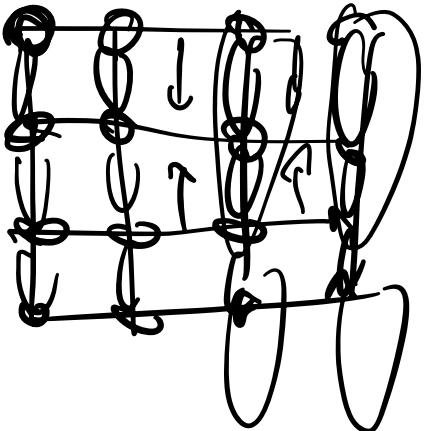
$$\text{ej: } f \rightarrow \tilde{f} = \frac{f}{b}$$

4) La aproximación de Kadanoff

para  $\mathbb{P}^D$ .

BS





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$$\hat{k}_g = b k_g \quad \hat{k}_x = \text{th}^{-1} [\text{th}[k_x]^b]$$

at final

$$F = b \text{th}^{-1} [\text{th}(bx)]^b$$

$$b = e^{\delta P} \sim (1 + \delta P)$$

$$\text{th}^{-1} [\text{th}[(1 + \delta P) x]^{(1 + \delta P)}]$$

$$\simeq (1 + \delta P) x + \delta P \left[ \sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{n} \ln^n \text{th}(x) \right] + O(\delta P^2)$$

$$K^{\text{ref}} = K + \frac{1}{2} \kappa$$

$$\frac{dK}{dt} = K + \frac{1}{2} [\sin(2\omega) \ln(\tan \omega)]$$

para  $D=2$

se muesta

$$\frac{dk}{dt} = (D-1)k + \frac{1}{2} [\sin(2\omega) \ln(\tan \omega)]$$

$$T = \frac{S}{k_B \omega}, \text{ para } T \text{ pequeño.}$$

$$\Rightarrow \frac{dT}{dt} = -\epsilon T + \frac{R_B}{2S} T^2 + O(T^3)$$

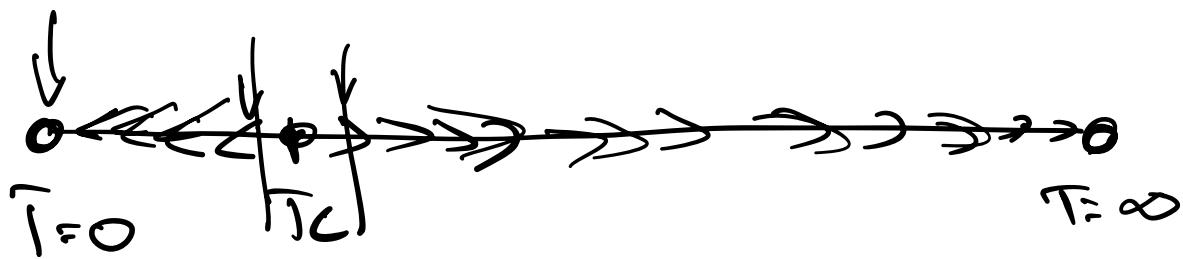
$$\epsilon = D-1$$

temperatura constante ( $\frac{dT}{dx} = 0$ )

$$T=0$$

$$T_c = \frac{2e\beta}{k_B}$$

$$\epsilon = D - 1$$



$$+ T = 0 + \delta T$$

$$D \frac{d\delta T}{dt} = -\epsilon \delta T \quad \epsilon > 0$$

$$+ T = T_c + \delta T$$

$$D \frac{d\delta T}{dt} = +\epsilon \delta T$$

$T_c$  es un punto fijo inestable.

$$\delta T(t) = \delta T(0) e^{\epsilon t}$$

}

$$g(\ell) = g(0) e^{-\ell}$$

$$g \sim \delta T^{-\nu}$$

$$\nu = \frac{1}{\epsilon} = \frac{1}{D-1}$$

5) La renormalización para la  
Teoría de campo.

$$S = \int d^D x \left[ \frac{\kappa}{2} (\vec{\nabla} \phi)^2 + \frac{t}{2} \phi^2 + \mu \phi^4 + \mu_6 \phi^6 + \dots \right]$$

Sabemos que:

$$\text{La acción } S = \int d^D x \frac{\kappa}{2} (\vec{\nabla} \phi)^2 \text{ es invariante}$$

de escala.

$$\bar{n} \rightarrow b \bar{n}' \quad \kappa \rightarrow \kappa'$$

$$\phi \rightarrow b^{\frac{r-2}{2}} \phi'$$

$\alpha$  about  $\neq 0$ ,  $a \neq 0$ ,  $m_0 \neq 0$  etc..

$$t \rightarrow b^{-2} t' \Rightarrow t' = b^2 t^{4-D}$$

$$\mu \rightarrow b^{-4+D} \mu' \Rightarrow \mu' = b^{4-D} \mu$$

$$m_0 \rightarrow b^{-6+2D} m_0' \quad m_0' = b^{\frac{2D}{2D+6}} m_0 \Rightarrow b^{\frac{2}{D+3}} m_0$$

$$u_g \rightarrow b^{-8+3D} u_g'$$

obs in  $D=4$   $E=0$ ,  $m_0 = u_g = \dots = 0$

$$\mu \neq 0$$

$$S = \int d^4x \left[ \frac{1}{2} (\nabla \phi)^2 + m \phi^4 \right]$$

in  $D=4$  es invariante de escala!

para  $D=3$ ,  $\mu > m$