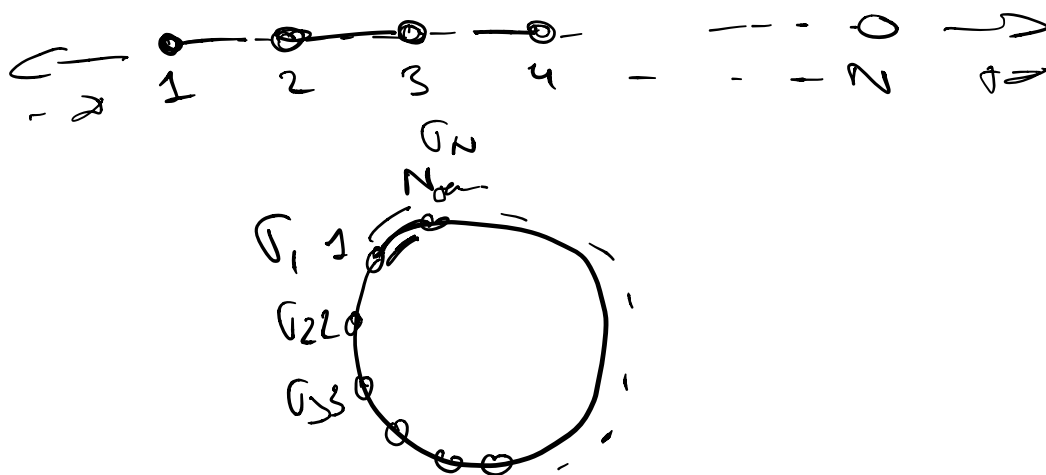


# 8) Modelo de Ising em 1-D

Cond. de Borda periódicas.



$$\underline{\sigma_{N+1} \equiv \sigma_1} \quad ; \quad \sigma_{N+2} \equiv \sigma_2$$

$$H = -J \sum_{i=1}^N \sigma_i \sigma_{i+1} - h \sum_{i=1}^N \sigma_i$$

$$Z = \sum_{\{\sigma_i\}} e^{-\beta H} \quad ; \quad K = \frac{J}{k_B T} = \beta J$$

$$\hat{h} = \frac{h}{k_B T}$$

$$Z = \sum_{\{\sigma\}} e^{K \sum_{i=1}^N \sigma_i \sigma_{i+1} + \frac{h}{2} \sum_{i=1}^N \sigma_i + \frac{h}{2} \sum_{i=1}^N \sigma_{i+1}}$$

$$= \sum_{\{\sigma\}} e^i \sum_j \left( K \sigma_i \sigma_{i+1} + \frac{h}{2} \sigma_i + \frac{h}{2} \sigma_{i+1} \right)$$

$\{\sigma\}$   
 $\sigma_{1 \pm 1}, \sigma_2 = \pm 1$   
 $\dots \sigma_N = \pm 1$

$$= \sum_{\substack{i_1=1,2 \\ i_2=1,2 \\ \dots \\ i_N=1,2}} T_{i_1, i_2} T_{i_2, i_3} T_{i_3, i_4} \dots T_{i_N, i_1}$$

matriz de transferencia

$$T_{11} = \frac{e^{K \sigma_i \sigma_{i+1} + \frac{h}{2} \sigma_i + \frac{h}{2} \sigma_{i+1}}}{\left. \begin{array}{l} \sigma_i = 1 \\ \sigma_{i+1} = 1 \end{array} \right\}}$$

$$T_{12} = \dots \left. \begin{array}{l} \sigma_i = 1 \\ \sigma_{i+1} = -1 \end{array} \right\}$$

$$T_{21} = \dots \left. \begin{array}{l} \sigma_i = -1 \\ \sigma_{i+1} = 1 \end{array} \right\}$$

$$T_{22} = \dots \left. \begin{array}{l} \sigma_i = \sigma_{i+1} = -1 \end{array} \right\}$$

$$T = \begin{pmatrix} e^{\kappa h} & e^{-\kappa} \\ e^{-\kappa} & e^{\kappa h} \end{pmatrix}$$

$$\mathcal{L} = \text{tr} \{ T^N \}$$

$$\sum_i T_{ii} = \text{tr} T$$

$$\exists \Theta \text{ s.t. } \Theta^t \Theta = \mathbb{I}_2 = \underline{\underline{I}} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$\Theta^t T \Theta = \begin{pmatrix} \epsilon_1 & 0 \\ 0 & \epsilon_2 \end{pmatrix} = \mathbb{D}$$

$$\Theta^t T^N \Theta = \underbrace{\Theta^t T}_{\hat{\Theta}^t} \cdot \underbrace{T}_{\hat{\Theta}^t} \cdot \underbrace{T}_{\hat{\Theta}^t} \cdot \Theta = \mathbb{D}^N$$

$$\Theta^t T^N \Theta = \mathbb{D}^N \quad \text{tr} \{ \Theta^t T^N \Theta \}$$

$$\text{tr} \{ \mathbb{D}^N \} = \text{tr} \left\{ T^N \underbrace{\Theta \Theta^t}_{\mathbb{I}_d} \right\} = \text{tr} \{ T^N \}$$

$$\mathbb{D}^N = \begin{pmatrix} \epsilon_1^N & 0 \\ 0 & \epsilon_2^N \end{pmatrix}$$

$$\text{tr} \left\{ \overline{T^N} \right\} = \epsilon_1^N + \epsilon_2^N$$

$$\epsilon_{1/2} = \frac{e^{\kappa} \text{ch}(\tilde{h}) \pm \left( e^{2\kappa} \text{sh}^2(\tilde{h}) + e^{-2\kappa} \right)^{1/2}}{2} \quad \epsilon_j.$$

$$\epsilon_1 > \epsilon_2 > 0 \quad \epsilon_j.$$

$$Z = \epsilon_1^N + \epsilon_2^N \quad \text{as } N \gg 1$$

$$= \epsilon_1^N \left( 1 + \left( \frac{\epsilon_2}{\epsilon_1} \right)^N \right) \approx \epsilon_1^N$$

$$\frac{\epsilon_2}{\epsilon_1} < 1$$

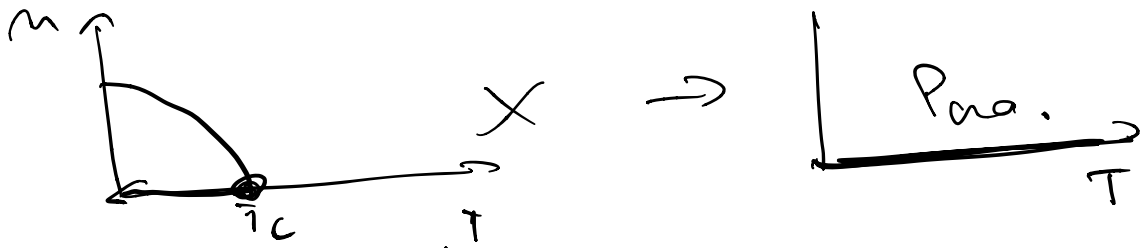
$$F = -k_B T \ln Z \approx -k_B T N \ln \epsilon_1$$

$$m = \frac{1}{N} \frac{\partial}{\partial h} \ln Z = \frac{\partial}{\partial h} \ln \epsilon_1$$

$$\Rightarrow m = \frac{\text{sinh}(\tilde{h})}{k_B T \sqrt{\text{sinh}^2(\tilde{h}) + e^{-4\kappa}}}$$

obs as  $\tilde{h} \rightarrow 0$ ,  $m \rightarrow 0 \quad \forall T.$

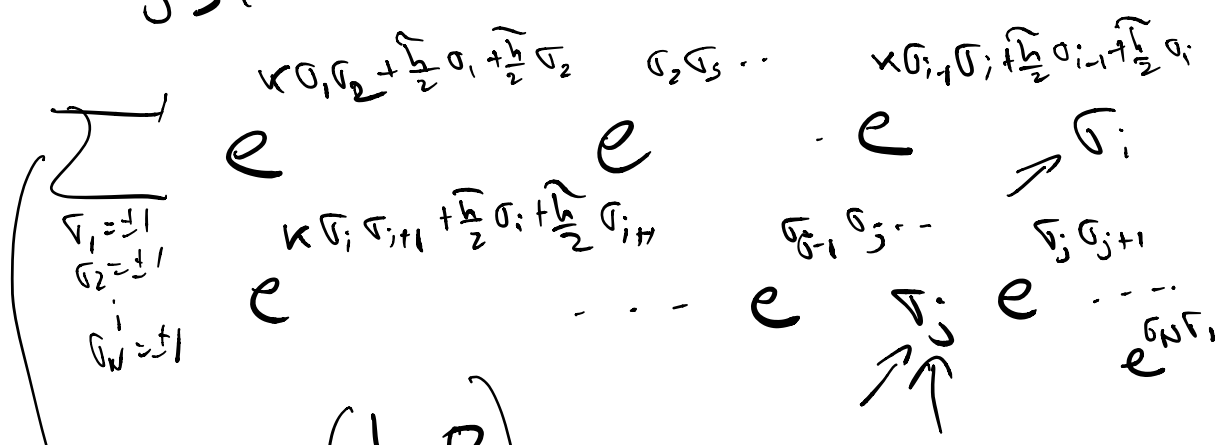
$\rightarrow$  no hay fase ferromagnética!



funciones de correlación: h=0

$$\langle \sigma_i \sigma_j \rangle = \frac{1}{Z} \sum_{\{\sigma\}} e^{-\beta H} \sigma_i \sigma_j$$

$j > i$



$$\Sigma = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$= \text{tr} \left\{ \underbrace{T \cdot T \cdot T \dots T}_i \cdot \underbrace{\Sigma}_{j-i} \cdot \underbrace{T \cdot T \cdot T \dots T}_{j-i} \cdot \underbrace{\Sigma}_{N-j} \cdot \underbrace{T \cdot T \cdot T \dots T}_{N-j} \right\}$$

$$\Rightarrow \langle \sigma_i \sigma_j \rangle = \frac{1}{Z} \text{tr} \left\{ T^i \Sigma T^{j-i} \Sigma T^{N-j} \right\}$$

$$T = \begin{pmatrix} e^k & e^{-k} \\ e^{-k} & e^k \end{pmatrix} : \Theta = \Theta^t = \begin{pmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} \end{pmatrix}$$

$$\Theta^t \Theta = \Theta \Theta \equiv \mathbb{I}_d$$

$$\Theta T \Theta = \begin{pmatrix} t_1 & 0 \\ 0 & t_2 \end{pmatrix}$$

$$\Theta T^n \Theta = \begin{pmatrix} t_1^n & 0 \\ 0 & t_2^n \end{pmatrix}$$

↓

$$T \cdot T \cdot T \cdot T$$

$$\Rightarrow \langle \sigma_i \sigma_j \rangle = \frac{1}{2} \left\{ \underbrace{D^i \Theta \Sigma \Theta D^j \Theta \Sigma \Theta D^i \Theta \Sigma \Theta D^j}_{2} \right\}$$

$$\Theta \Sigma \Theta = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$\Rightarrow \langle \sigma_i \sigma_j \rangle = \frac{1}{2} \left\{ \begin{pmatrix} t_1^{i_0} & 0 \\ 0 & t_2^{i_0} \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} t_1^{j_0} & 0 \\ 0 & t_2^{j_0} \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \right.$$

$$\dots \begin{pmatrix} t_1^{i_{2j}} & 0 \\ 0 & t_2^{i_{2j}} \end{pmatrix} \dots$$

$$\Rightarrow \langle \sigma_i \sigma_j \rangle = \frac{\epsilon_1^{N-j+1} \epsilon_2^{j-i} + \epsilon_2^{N-j+1} \epsilon_1^{j-i}}{\epsilon_1^N + \epsilon_2^N}$$

$j-i$  fija,  $N \rightarrow \infty$

$$\langle \rangle \sim \frac{\cancel{\epsilon_1^N} (\epsilon_2^{j-i} \epsilon_1^{i-j} + \dots)}{\cancel{\epsilon_1^N} (1 + \dots)}$$

$$\Rightarrow \langle \sigma_i \sigma_j \rangle = \left( \frac{\epsilon_2}{\epsilon_1} \right)^{j-i} = [\text{tgh}(k)]^{j-i}$$

$$= e^{-\frac{j-i}{\xi}}$$

$$\text{tgh}(k) \equiv e^{-\frac{1}{\xi}}$$

longitud de correlación:

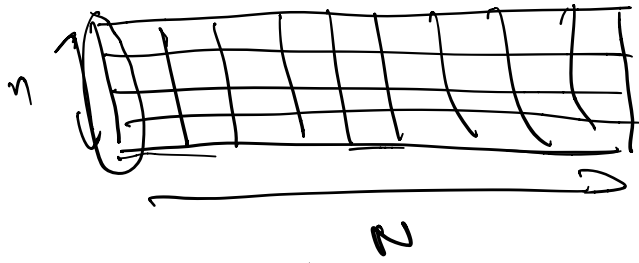
$$\xi = \frac{1}{\ln(\text{chgh}(k))} = \frac{1}{\ln(\text{chgh}(\frac{J}{k_B T}))}$$

obs ni  $T \rightarrow \infty$ ,  $\xi \rightarrow 0$

ni  $T \rightarrow 0$ ,  $\xi \rightarrow \infty$

1999, Onsager.

$\rightarrow$  Sol. para 2  $\rightarrow$



$$Z = \text{tr}(T^N)$$

n spins.  $\rightarrow 2^n$  posibilidades

$$T_{2^n \times 2^n} \rightarrow \epsilon_1, \epsilon_2, \dots, \epsilon_{2^n}$$

$\uparrow$              $\uparrow$   
 (Two arrows point from the ground line to the first two eigenvalues,  $\epsilon_1$  and  $\epsilon_2$ )

$$n \sim 20$$

$$n \sim 6, 8, 10$$

Kerson Huang.

## II Del modelo de Zring a la "acción" de Ginzburg-Landau.

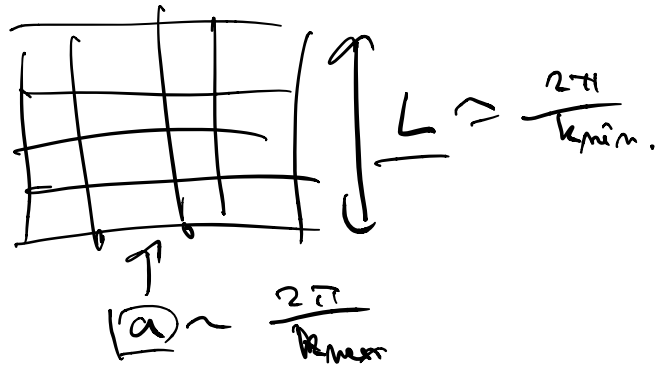
$$\int dk f(k) \quad \text{problema si } k \rightarrow \infty$$

cut-off  $k_{\text{max}} \rightarrow$  cut-off ultra-violetta.

$$\int \frac{dk}{k^n}$$

cut-off  $k_{\text{min}} = \Lambda$  cut-off Infra-rojo





$$\frac{a}{L} \rightarrow 0$$

$$\underline{L} \rightarrow \infty$$

1) transformación de Hubbard-Stratonovich

$$H = -\frac{1}{2} \sum_{\substack{\uparrow \\ i,j}} J_{ij} \underbrace{\sigma_i \sigma_j}_{\substack{\uparrow \\ \sigma_i \sigma_j}} \quad (\mathcal{J})_{N \times N}$$

$J_{ij}$  es una matriz simétrica.

$$M = M_S + M_A$$

$$M = \frac{M + M^T}{2} + \frac{M - M^T}{2}$$

$J_{ij}$  definida positiva (todos sus autovalores son  $> 0$ )

$P \rightarrow \lambda_1, \lambda_2, \dots, \lambda_N$  autovalores.

$$\tilde{M} = P + \alpha I$$

$\tilde{M} \rightarrow \lambda_1 + \alpha, \lambda_2 + \alpha, \dots, \lambda_N + \alpha$   
 $\alpha \neq \lambda_i + \alpha > 0$

obs  $\tilde{J} \Rightarrow \tilde{J} + \alpha I$

$$\sum_{i,j} \alpha (\beta_{ij})_{ij} \sigma_i \sigma_j = \underline{\underline{N\alpha}}$$

$\sigma_i^2 = 1$

$(J)_{N \times N}$  matriz de coherencia.

$\alpha \quad |i-j| \rightarrow \infty \quad J_{ij} \rightarrow 0$   
 $\uparrow$   
 distancia entre  $i$  y  $j$

$$f = -\frac{1}{2} \sigma^t J \sigma$$

$$\sigma^t J \sigma = (\sigma_1, \sigma_2, \dots, \sigma_N) J \begin{pmatrix} \sigma_1 \\ \vdots \\ \sigma_N \end{pmatrix}$$

$$Z = \sum_{\{0\}} e^{\frac{1}{2} \sigma^T \{0\}}$$

ej:  $\det A_{n \times n} > 0$  min. y real.

$$\begin{aligned}
 I &= \int_{-\infty}^{\infty} \prod_{i=1}^n dx_i e^{-\frac{1}{2} \sum_{i,j} x_i A_{ij} x_j + \sum_{i=1}^n x_i Y_i} \\
 &= \int_{-\infty}^{\infty} dx_1 dx_2 \dots dx_n e^{-\frac{1}{2} \vec{x}^T A \vec{x} + \vec{x}^T \cdot \vec{Y}} \\
 &= \int_{-\infty}^{\infty} \frac{d\vec{x}}{dx_1 \dots dx_n} e^{-\frac{1}{2} \sum_{i,j} Y_i (A^{-1})_{ij} Y_j} \\
 &= \frac{(2\pi)^n}{\det(A)} e^{-\frac{1}{2} \sum_{i,j} Y_i (A^{-1})_{ij} Y_j}
 \end{aligned}$$

Dem  $\vec{x} = \vec{x} + A^{-1} \vec{Y} \rightarrow x_i = \tilde{x}_i + \sum_j (A^{-1})_{ij} Y_j$

$$d\vec{x} = d\tilde{x}$$

$$\begin{aligned}
 I &= \int_{\tilde{x}} e^{-\frac{1}{2} \tilde{x}^T A \tilde{x}} \times e^{-\frac{1}{2} \vec{Y}^T A^{-1} \vec{Y}}
 \end{aligned}$$

$$\exists \theta \text{ s.t. } \theta^T \theta = I_n$$

$$\text{t.q. } \theta^T A \theta = D = \begin{pmatrix} \alpha_1 & & & 0 \\ & \alpha_2 & & \\ & & \ddots & \\ 0 & & & \alpha_n \end{pmatrix}$$

$$\mathbb{R}^n = \theta \mathbb{R}^n$$

$$\det \theta = 1$$

$$d\mathbb{R}^n = d\mathbb{R}^n$$

at last!

$$\int_{-\infty}^{\infty} dx e^{-\frac{\alpha}{2} x^2} = \sqrt{\frac{2\pi}{\alpha}} \quad \alpha > 0$$

$$\det(A) = \alpha_1 \alpha_2 \dots \alpha_n$$

$$\beta \mathcal{Z} \leftrightarrow A^{-1}$$

$$y_i \leftrightarrow \sigma_i$$

$$x_i \leftrightarrow \phi_i$$

$$\rho_{\mathbb{R}^n}^{\beta \mathcal{Z}} = (2\pi)^{-N/2} [\det(\beta \mathcal{Z})]^{-1/2} e^{-\frac{1}{2\beta} \bar{\phi}^T \beta \phi + \sum_{i=1}^n \sigma_i \phi_i} \times \int \prod_{i=1}^n d\phi_i$$

$$= \left( (2\pi)^{-N/2} [\text{Det}(\beta S)]^{-1/2} \int \prod_i d\phi_i e^{-\frac{1}{2\beta} \sum_{ij} \phi_i (S^{-1})_{ij} \phi_j + \sum_i \sigma_i \phi_i} \right)$$

Però  $Z = \sum_{\{\sigma\}} e^{\frac{\beta}{2} \vec{\sigma}^T S \vec{\sigma}}$

$$\Rightarrow Z = (2\pi)^{-N/2} [\text{Det}(\beta S)]^{-1/2} \int \prod_i d\phi_i e^{-\frac{1}{2\beta} \vec{\phi}^T S^{-1} \vec{\phi}} \underbrace{\sum_{\{\sigma\}} e^{\sum_i \sigma_i \phi_i}}$$

$$\sum_{\{\sigma\}} e^{\sum_i \sigma_i \phi_i} = \left( \sum_{\sigma=\pm 1} e^{\sigma \phi_1} \right) \left( \sum_{\sigma=\pm 1} e^{\sigma \phi_2} \right) \dots$$

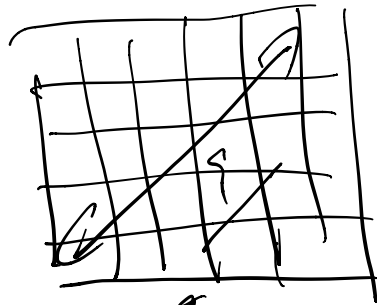
$2 \text{ch}(\phi_1) \quad 2 \text{ch}(\phi_2) \dots$

$$\rightarrow = 2^N \prod_i \text{ch}(\phi_i) = 2^N e^{\sum_i \ln[\text{ch}(\phi_i)]}$$

$$\Rightarrow Z = (2\pi)^{-N/2} 2^N [\text{Det}(\beta S)]^{-1/2} \int \prod_i d\phi_i e^{-\frac{1}{2\beta} \vec{\phi}^T S^{-1} \vec{\phi} - V(\vec{\phi})}$$

$$\text{con } U(\Phi) = - \sum_{i=1}^N \ln[\text{ch}(\Phi_i)]$$

## 2) El límite continuo



$\xi > a$   
nitrato

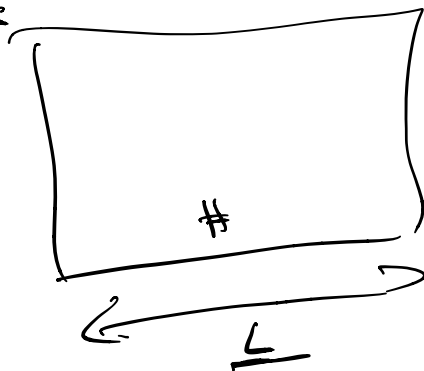
$a$

$$\langle \sigma_i \sigma_j \rangle \sim \frac{C}{|i-j|^\alpha}$$

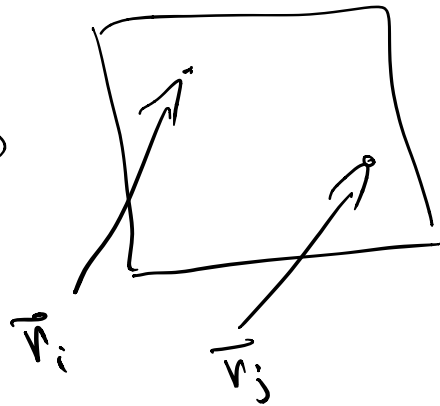
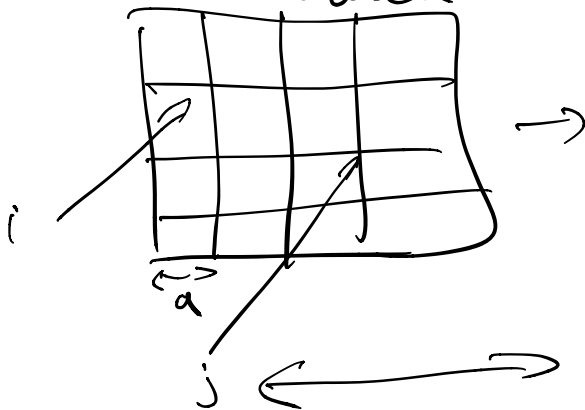
$$N \rightarrow \infty$$

$$a \rightarrow \infty$$

$$L \gg a$$



dimensionalidad  $d$



$$\sum_i \rightarrow \int \frac{d^d x}{|a_d|}$$

$$C \delta_{i,j} \rightarrow a^d \delta(\vec{r}_i - \vec{r}_j)$$

$$\sum_j \delta_{i,j} = 1 \rightarrow \frac{1}{a^d} \int d\vec{r}_j a^d \delta(\vec{r}_i - \vec{r}_j) = 1$$

$J_{ij}$  depende de  $|i-j|$

$$\hookrightarrow J(|\vec{r}_i - \vec{r}_j|)$$

$$J_{ij}^{-1} \rightarrow J^{-1}(|\vec{r}_i - \vec{r}_j|)$$

$$\sum_j J_{ij} J_{jk}^{-1} = \delta_{i,k}$$

$$\frac{1}{a^d} \int d\vec{r}_j J(|\vec{r}_i - \vec{r}_j|) J^{-1}(|\vec{r}_j - \vec{r}_k|) = a^d \delta(|\vec{r}_i - \vec{r}_k|) \quad (*)$$

$$\phi_i \rightarrow \phi(\vec{r}_i)$$

definimos las transformadas de Fourier

$$\hat{\phi}(\vec{k}) = \int d\vec{r} e^{i\vec{k} \cdot \vec{r}} \phi(\vec{r}) \quad (\text{T.f.})$$

$$\phi(\vec{r}) = \frac{1}{(2\pi)^d} \int d^d \vec{k} e^{-i\vec{k} \cdot \vec{r}} \hat{\phi}(\vec{k})$$

obs  $\phi(\vec{r})$  is real  $\Rightarrow \hat{\phi}^*(\vec{k}) = \hat{\phi}(\vec{-k})$

$$\hat{J}(\vec{k}) = \int d^d \vec{r} e^{i\vec{k} \cdot \vec{r}} J(\vec{r})$$

$$J(\vec{r}) = \frac{1}{(2\pi)^d} \int d^d \vec{k} e^{-i\vec{k} \cdot \vec{r}} \hat{J}(\vec{k})$$

$$\hat{J}^{-1}(\vec{k}) = \int d^d \vec{r} e^{i\vec{k} \cdot \vec{r}} J^{-1}(\vec{r})$$

$$\int d^d \vec{r} e^{i(\vec{k} + \vec{k}') \cdot \vec{r}} = (2\pi)^d \delta(\vec{k} + \vec{k}')$$

$$\int d^d \vec{r} e^{i\vec{k} \cdot (\vec{r} - \vec{r}')} = (2\pi)^d \delta(\vec{r} - \vec{r}')$$

g.g.:

$$\hat{J}^{-1}(\vec{k}) = \frac{a^{2d}}{\hat{J}(\vec{k})}$$

$$\sum_{i,j} \phi_i \hat{J}_{ij}^{-1} \phi_j \rightarrow \frac{1}{a^{2d}} \int d^d \vec{r} d^d \vec{r}' \phi(\vec{r}) \hat{J}^{-1}(\vec{r} - \vec{r}') \phi(\vec{r}')$$



$$\begin{aligned} \xrightarrow{\text{T.f.}} & \frac{1}{a^{2d}} \frac{1}{(2\pi)^d} \int d\vec{k} |\hat{\phi}(\vec{k})|^2 \hat{J}^{-1}(\vec{k}) \\ & \left( \hat{\phi}(\vec{k}) = \hat{\phi}^*(\vec{k}) \right) \\ & = \frac{1}{(2\pi)^d} \int d\vec{k} \frac{|\hat{\phi}(\vec{k})|^2}{\hat{J}(\vec{k})} \end{aligned}$$


---

→ distancias grandes  
 →  $|\vec{k}|$  pequeños.

$$\int e^{ikx} f(x) \rightarrow \hat{f}(k) \xrightarrow{k \rightarrow 0} 0$$

$$x \rightarrow x + \delta x \quad \delta x \ll \lambda$$

$$f(x + \delta x) \approx f(x)$$

$$e^{ikx + ik\delta x}$$

$$e^{ik\delta x} \neq 1$$

$$\int_{k_{\min}}^{k_{\max}} dk e^{ikx} \approx \delta k \rightarrow \frac{1}{x}$$

$$\hat{J}(k) = \int d^n \vec{n} e^{i\vec{k} \cdot \vec{n}} \underline{J(|\vec{n}|)}$$

$$\hat{J}(k) = \bar{J}(0) + \sum_{\alpha=1, \dots, d} k^\alpha \left. \frac{\partial \bar{J}}{\partial k^\alpha} \right|_{k=0} + \frac{1}{2} \sum_{\alpha, \beta} k^\alpha k^\beta \left. \frac{\partial^2 \bar{J}}{\partial k^\alpha \partial k^\beta} \right|_{k=0} + \mathcal{O}(k^3) + \mathcal{O}(k^4)$$

$$\int d^n \vec{n} n^\alpha \bar{J}(|\vec{n}|)$$

$\lim_{n \rightarrow -n} n^\alpha \rightarrow -n^\alpha$

$$\int d^n \vec{n} n^\alpha n^\beta \bar{J}(|\vec{n}|) \quad n^\alpha \neq n^\beta = 0$$

$$\hat{J}(k) = \hat{J}(0) + \frac{1}{2} \sum_{\alpha} k^{\alpha 2} \left. \frac{\partial^2 \bar{J}}{\partial k^{\alpha 2}} \right|_{k=0} + \dots$$

$$= \frac{1}{2} \int d^n \vec{n} |\vec{n}|^2 \bar{J}(|\vec{n}|)$$

$l$ : radio de alcance de la matriz de conectividad.

$$\text{Def } l^2 \equiv \frac{\int d^n \vec{r} \, |\vec{r}|^2 \mathcal{Z}(\vec{r})}{\int d^n \vec{r} \, \mathcal{Z}(\vec{r})}$$

$$\Rightarrow \hat{\mathcal{Z}}(\vec{r}) = \hat{\mathcal{Z}}(0) [1 - l^2 \vec{r}^2 + \dots]$$

$$\frac{1}{\hat{\mathcal{Z}}(\vec{r})} \sim \frac{1}{\hat{\mathcal{Z}}(0)} [1 + l^2 \vec{r}^2 + \dots]$$

$$\Rightarrow \frac{1}{2\beta} \vec{\phi}^T \hat{\mathcal{Z}}^{-1} \vec{\phi} \rightarrow \frac{1}{2\beta \hat{\mathcal{Z}}(0)} \frac{1}{(2\pi)^d} \int d^n \vec{k} \dots$$

$$\dots |\hat{\phi}(\vec{r})|^2 [1 + l^2 \vec{r}^2 + \dots]$$

$$\left( = \frac{1}{2\beta \hat{\mathcal{Z}}(0)} \int d^n \vec{r} [\phi^2(\vec{r}) + l^2 |\nabla \phi|^2 + \dots] \right)$$

$$\sum_i \ln \text{ch } \phi_i \rightarrow \frac{1}{a^d} \int d^n \vec{r} \ln [\text{ch}(\phi)]$$

$$\ln [\text{ch}(\phi)] \sim \ln \left[ 1 + \frac{\phi^2}{2} + \frac{\phi^4}{4!} + \dots \right] = \frac{\phi^2}{2} - \frac{\phi^4}{12}$$

...

$$e^{-\frac{1}{2\beta} \vec{\phi}^T \hat{S} \vec{\phi} - U(\vec{\phi})} = e^{-\frac{1}{2\beta \hat{S}(0)} \int d^n \vec{r} [\phi(\vec{r})^2 + \lambda^2 |\nabla \phi|^2 - \frac{1}{a_2} (\frac{\phi^2}{2} - \frac{\phi^4}{4})]}$$

$$Z = (2\pi)^{-N/2} (\text{Det}(\beta \hat{S}))^{-1/2} Z^N$$

$$\int \mathcal{D}(\phi) = \lim_{N \rightarrow \infty} \int \prod_i \phi_i$$

acción de Ginzburg-Landau.

$$S(\phi) = \int d^n \vec{r} \left[ \frac{c}{2} (\nabla \phi)^2 + \frac{a_2}{2} \phi^2 + a_4 \phi^4 + \dots \right]$$

$$c = \frac{\lambda^2}{\beta \hat{S}(0)} ; a_2 = \frac{k_B T}{\hat{S}(0)} - \frac{1}{a_1} , a_4 = \frac{1}{12 a_1^2}$$

$$Z = \int \mathcal{D}\phi e^{-S(\phi)}$$

$$\begin{pmatrix} - \\ - \\ 1 \\ - \\ 1 \end{pmatrix} \rightarrow \begin{pmatrix} \phi \\ 1 \\ 1 \end{pmatrix}$$

$$(\nabla \phi)^2 \leftrightarrow \partial_m \phi \partial_m \phi$$

$$a_2 \leftrightarrow m^2, a_4 \leftrightarrow \lambda$$

$$e^{\frac{i}{\hbar} S} \quad e^{\frac{i}{\hbar} \int_{t'}^t \mathcal{L} dt}$$

relación de Wick  $t = iz$

$$\partial_t^2 \rightarrow -\partial_z^2$$

$$\partial_{t'}^2 \rightarrow -\partial_{t'}^2 \rightarrow \Delta$$

$$\begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix} \rightarrow \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix}$$

$$U \rightarrow e^{\frac{i}{\hbar} H(t-t')}$$

$$e^{\frac{i}{\hbar} H \Delta t} \quad e^{\frac{i}{\hbar} H \Delta t}$$

$$\int \mathcal{D}q \mathcal{D}p$$

$$e^{-\beta H} \rightarrow e^{-\Delta \beta H} \quad e^{-\Delta \beta H}$$

$$\langle \bar{S}_i \bar{S}_j \rangle$$

$$\bar{S}_i^2, \bar{S}_j^2 = \hbar S(S+1)$$

$$\left( \overline{S_i - S_j} \right)^2$$